

Overview to mathematical analysis for fractional diffusion equations – new mathematical aspects motivated by industrial collaboration

Junichi Nakagawa, Kenichi Sakamoto and Masahiro Yamamoto

Received on March 31, 2010

Abstract. The mathematics turns out to be useful for creation of innovations in the industry, and the mathematical knowledge and thinking manners are used effectively for that purpose. However, this is only one aspect of the industrial mathematics where various existing mathematical knowledge are applied for solving required subjects from industry. On the other hand, one can see the opposite direction; Pursuit of industrial purposes inspires to create new fields of mathematics by motivating and activating existing researches. is an important aspect of the industrial mathematics because it does not only give tools for solving concrete problems, but also enriches the existing branches of mathematics. In this article, as such a possible example, we discuss a fractional diffusion equation which has been studied already comprehensively from the theoretical interests, but the researches are expanded as a mathematical topic in view of the industrial applications.

Keywords. mathematics motivated by industrial mathematics, fractional diffusion equation, fractional calculus, well-posedness, qualitative properties

1. INTRODUCTION

The diffusion of contaminants under the ground is important and from the environmental viewpoint, better simulations and predictions of the density of the contaminant over time should be done. Moreover the real size is over a few kilometers, while one can execute only laboratory experiments with meter sizes (see Figure 1).

As classical model equation, one can use a diffusion convection equation:

$$\rho(x) \frac{\partial u}{\partial t}(x, t) = \operatorname{div}(p(x) \nabla u(x, t)) + b(x) \cdot \nabla u(x, t),$$

where $u(x, t)$ denotes the density at time t and the location x . In 1992, Adams and Gelhar [1] pointed that field data show anomalous diffusion in heterogeneous aquifer which can not be interpreted by the classical convection-diffusion equation (see Figure 2). Since [1], there are trials for better modelling and we can refer to Berkowitz, Scher and Silliman [6], Y. Hatano and N. Hatano [18]. See also Berkowitz, Cortis, Dentz and Scher [5], Xiaong, G. Huang and Q. Huang [52]. The diffusion is observed to be slower than the prediction on the basis of the classical convection-diffusion equation, and such anomalous diffusion is called "slow diffusion". We refer especially to Y. Hatano and N. Hatano [18] where the continuous-time random walk is discussed. In the soil, one has to take into consideration the porosity and the heterogeneity of the medium, and by the microscopic level, one can conclude that the classical random walk model may not be suitable in view of the heterogeneity. The continuous-time random walk is a microscopic

model for the anomalous diffusion, and by an argument similar to the derivation of the classical diffusion equation from the random walk, one can derive fractional diffusion models (e.g., Metzler and Klafter [34], (pp.14-18), Sokolov, Klafter and Blumen [50]).

The fractional diffusion equation can be described as follows. Let $0 < \alpha < 1$ throughout this paper. We consider

$$\partial_t^\alpha u(x, t) = (Lu)(x, t) + F(x, t), \quad x \in \Omega, t \in (0, T), \quad (1.1)$$

where $\Omega \subset R^n$ is a bounded domain with smooth boundary $\partial\Omega$, ∂_t^α denotes the Caputo fractional derivative with respect to t and is defined by

$$\partial_t^\alpha u(x, t) = \frac{1}{\Gamma(1-\alpha)} \int_0^t (t-\tau)^{-\alpha} \frac{\partial u}{\partial \tau}(x, \tau) d\tau$$

for x -dependent function $u(x, t)$ and

$$D_t^\alpha g(t) = \frac{1}{\Gamma(1-\alpha)} \int_0^t (t-\tau)^{-\alpha} \frac{dg}{d\tau}(\tau) d\tau$$

for x -independent function $g(t)$ (e.g., Podlubny [41]), Γ is the Gamma function and the operator L is a symmetric uniformly elliptic operator:

$$(Lu)(x) = \sum_{i,j=1}^n \frac{\partial}{\partial x_i} \left(a_{ij}(x) \frac{\partial u}{\partial x_j}(x) \right) + c(x)u(x), \quad x \in \Omega,$$

where $a_{ij} = a_{ji}$, $\in C^1(\bar{\Omega})$, $c \in C(\bar{\Omega})$, ≤ 0 on $\bar{\Omega}$ and we assume that there exists a constant $\mu > 0$ such that

$$\sum_{i,j=1}^n a_{ij}(x) \zeta_i \zeta_j \geq \mu \sum_{i=1}^n \zeta_i^2$$

for all $x \in \bar{\Omega}$ and $\zeta_1, \dots, \zeta_n \in R$. Moreover F is a given function in $\Omega \times [0, T]$ and $T > 0$ is a fixed value.

The fractional diffusion equation needs independent mathematical researches, even though one can discuss similarly to the classical convection-diffusion equation. One has to take into consideration that some properties for the natural number order derivatives fail for fractional order derivatives: For example, the derivative of the product of two functions and the sequential derivative do not hold.

We note that

$$\lim_{\alpha \rightarrow 1} D_t^\alpha g(t) = \frac{dg}{dt}(t), \quad 0 \leq t \leq T$$

for $g \in C^2[0, T]$. In fact, the integration by parts yields

$$D_t^\alpha g(t) = \frac{1}{\Gamma(2-\alpha)} \left(g'(0)t^{1-\alpha} + \int_0^t (t-s)^{1-\alpha} g''(s) ds \right) \rightarrow g'(t)$$

as $\alpha \rightarrow 1$ for arbitrary $t \in [0, T]$.

This means that the Caputo derivative of order $\alpha \in (0, 1)$ has an extended sense of the first-order derivative.

As theoretical backgrounds for e.g., better simulation requested for the environmental or possible industrial applications, one can apply mathematical results which have been already gained. However, in view of the applications, further mathematical researches may be necessary, which is quite a strong motivation for mathematicians and may open new aspects of the vast field of the fractional differential equation. That is, for better applications, mathematicians should sometimes modify the existing theories and even create and develop new branches in mathematics. This is bilaterally meaningful collaboration between mathematics and industry. We expect that the fractional diffusion equation may be such a topic. In this article, we intend a compact overview to such aspects concerning the fractional differential equations and present results which have been proved by the authors' group and their colleagues. As for more complete descriptions and the proofs, we refer to the original papers, e.g., Cheng, Nakagawa, Yamamoto and Yamazaki [7], Sakamoto [45] and Sakamoto and Yamamoto [46], and we omit.

The article is composed of 5 sections. In section 2, we discuss some specific aspects of fractional calculus and in section 3, we choose topics on ordinary fractional differential equations. In section 4, we will present results on the well-posedness of initial/boundary value problems for fractional diffusion equations to show qualitative properties which interpret the character as slow diffusion and in section 5 we discuss more properties related with inverse problems.

2. FRACTIONAL CALCULUS

For a function $g \in C^1[0, T]$, we recall

$$D_t^\alpha g(t) = \frac{1}{\Gamma(1-\alpha)} \int_0^t (t-\tau)^{-\alpha} \frac{d}{d\tau} g(\tau) d\tau$$

for $0 < \alpha < 1$. By the definition, for example, we can calculate:

$$D_t^\alpha t^\gamma = \frac{\Gamma(\gamma+1)}{\Gamma(\gamma+1-\alpha)} t^{\gamma-\alpha}, \quad \gamma > 0. \quad (2.1)$$

The Caputo derivative is defined by the integral and so is not a local operation, and several properties for the usual calculus do not hold.

First we note that we have no useful formula for the derivative of product of two functions:

$$D_t^\alpha (fg) \neq (D_t^\alpha f)g + fD_t^\alpha g,$$

and accordingly we have no useful formula for the integration by parts (e.g., [41]).

Moreover we have no usual properties for sequential derivatives in general:

$$D_t^\alpha D_t^\beta \neq D_t^{\alpha+\beta}$$

even if $0 < \alpha, \beta < 1$ and $\alpha + \beta < 1$.

In fact, let $0 < \alpha < \frac{1}{2}$. By (2.1), we have $D_t^\alpha t^\alpha = \Gamma(\alpha+1)$, and $D_t^\alpha (D_t^\alpha t^\alpha) = 0$, but

$$D_t^{2\alpha} t^\alpha = \frac{\Gamma(1+\alpha)}{\Gamma(1-\alpha)} t^{-\alpha},$$

that is, $D_t^\alpha (D_t^\alpha t^\alpha) \neq D_t^{2\alpha} t^\alpha$. On the other hand, we note by (2.1) that

$$D_t^\beta (D_t^\alpha t^\gamma) = D_t^{\alpha+\beta} t^\gamma$$

if $\gamma - \alpha > 0$. More generally we can prove

Proposition 2.1

Let $f \in C^2[0, T]$ and let $0 < \alpha, \beta < 1$, $\alpha + \beta < 1$. Then

$$D_t^\beta (D_t^\alpha f)(t) = D_t^{\alpha+\beta} f(t), \quad 0 \leq t \leq T.$$

Proof. We have

$$\begin{aligned} D_t^\alpha f(t) &= \frac{1}{\Gamma(1-\alpha)} \int_0^t (t-s)^{-\alpha} f'(s) ds \\ &= \frac{1}{\Gamma(1-\alpha)} \int_0^t (t-s)^{-\alpha} (f'(s) - f'(t)) ds \\ &\quad + \frac{1}{(1-\alpha)\Gamma(1-\alpha)} t^{1-\alpha} f'(t). \end{aligned}$$

Then

$$\begin{aligned} (D_t^\alpha f)'(t) &= \frac{1}{\Gamma(1-\alpha)} \int_0^t \partial_t ((t-s)^{-\alpha}) \\ &\quad \times (f'(s) - f'(t)) ds \\ &\quad + \frac{1}{\Gamma(1-\alpha)} \left(- \int_0^t (t-s)^{-\alpha} f''(t) ds \right. \\ &\quad \left. + \frac{1}{1-\alpha} f''(t) t^{1-\alpha} + t^{-\alpha} f'(t) \right). \end{aligned}$$

Since $\partial_t((t-s)^{-\alpha}) = -\partial_s((t-s)^{-\alpha})$, we have

$$\begin{aligned} & \int_0^t \partial_t((t-s)^{-\alpha})(f'(s) - f'(t))ds \\ &= - \int_0^t \partial_s((t-s)^{-\alpha})(f'(s) - f'(t))ds \\ &= - \left[(t-s)^{-\alpha}(f'(s) - f'(t)) \right]_{s=0}^{s=t} \\ & \quad + \int_0^t (t-s)^{-\alpha} f''(s)ds \\ &= t^{-\alpha}(f'(0) - f'(t)) \\ & \quad + \int_0^t (t-s)^{-\alpha} f''(s)ds. \end{aligned}$$

Hence we have

$$(D_t^\alpha f)'(t) = \frac{1}{\Gamma(1-\alpha)} \left(\int_0^t (t-s)^{-\alpha} f''(s)ds + t^{-\alpha} f'(0) \right).$$

Therefore

$$\begin{aligned} (D_t^\beta D_t^\alpha f)(t) &= \frac{1}{\Gamma(1-\beta)} \int_0^t (t-s)^{-\beta} (D_s^\alpha f(s))' ds \\ &= \frac{1}{\Gamma(1-\alpha)\Gamma(1-\beta)} \left(\int_0^t (t-s)^{-\beta} \right. \\ & \quad \times \left. \left(\int_0^s (s-\xi)^{-\alpha} f''(\xi) d\xi \right) ds \right. \\ & \quad \left. + \int_0^t s^{-\alpha} (t-s)^{-\beta} ds f'(0) \right). \end{aligned}$$

Noting $\int_0^t \left(\int_0^s d\xi \right) ds = \int_0^t \left(\int_\xi^t ds \right) d\xi$ and

$$\begin{aligned} & \frac{1}{\Gamma(1-\alpha)\Gamma(1-\beta)} \int_0^t \left(\int_\xi^t (t-s)^{-\beta} \right. \\ & \quad \left. \times (s-\xi)^{-\alpha} ds \right) f''(\xi) d\xi \\ &= \frac{1}{\Gamma(2-\alpha-\beta)} \int_0^t (t-\xi)^{1-\alpha-\beta} f''(\xi) d\xi, \end{aligned}$$

by integration by parts, we have

$$\begin{aligned} & \frac{1}{\Gamma(1-\alpha)\Gamma(1-\beta)} \int_0^t \left(\int_\xi^t (t-s)^{-\beta} \right. \\ & \quad \left. \times (s-\xi)^{-\alpha} ds \right) f''(\xi) d\xi \\ &= \frac{1}{\Gamma(2-\alpha-\beta)} \left\{ \left[f'(\xi)(t-\xi)^{1-\alpha-\beta} \right]_{\xi=0}^{\xi=t} \right. \\ & \quad \left. + (1-\alpha-\beta) \int_0^t (t-\xi)^{-\alpha-\beta} f'(\xi) d\xi \right\}, \end{aligned}$$

that is, $(D_t^\beta D_t^\alpha f)(t) = D_t^{\alpha+\beta} f(t)$. The roof of the proposition is completed.

Moreover in Luchko [25], the following is proved.

Proposition 2.2

Let $g \in C^1[0, T]$ attain the maximum at $t = t_0 \in (0, T]$. Then

$$(D_t^\alpha g)(t_0) \geq 0.$$

On the other hand, we can not determine the local behaviour of g near $t = t_0$ by $D_t^\alpha g(t_0)$ because D_t^α is not a local operation.

As for further detailed account of fractional calculus, see Kilbas, Srivastava and Trujillo [20], Miller and Ross [35], Oldham and Spanier [38], Podlubny [41], Samko, Kilbas and Marichev [47].

3. ORDINARY FRACTIONAL DIFFERENTIAL EQUATIONS

For discussions about the fractional diffusion equation, the ordinary fractional differential equation is useful and is an independent important topic. We consider

$$D_t^\alpha u(t) = F(u, t), \quad t > 0, \quad u(0) = a. \quad (3.1)$$

Here $a \in \mathbb{R}$ and F is a given function. First let $F(u, t) = \lambda u + f(t)$, where λ is a constant:

$$\begin{aligned} D_t^\alpha u(t) &= \lambda u(t) + f(t), \quad t > 0, \\ u(0) &= a. \end{aligned} \quad (3.2)$$

In Gorenflo and Mainardi [15], Gorenflo and Rutman [17] (also see pp.140-141 in [20]), it is proved that there exists a unique solution to (3.2) and

$$\begin{aligned} u(t) &= aE_{\alpha,1}(\lambda t^\alpha) \\ & \quad + \int_0^t (t-s)^{\alpha-1} E_{\alpha,\alpha}(\lambda(t-s)^\alpha) f(s) ds \end{aligned} \quad (3.3)$$

for $t > 0$. Here $E_{\alpha,\beta}(t)$, $\alpha, \beta > 0$, is the Mittag-Leffler function:

$$E_{\alpha,\beta}(z) = \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(\alpha k + \beta)}$$

(see e.g., [41]) and is an entire function. Noting that

$$E_{1,1}(z) = e^z,$$

we see that (3.3) is $u(t) = ae^{\lambda t}$ for $\alpha = 1$ and $f \equiv 0$.

As for the unique existence of solution to (3.1), we can modify arguments in Chapter 3 in Kilbas, Srivastava and Trujillo [20] for example. See also Gorenflo and Mainardi [16].

Similarly to the ordinary differential equation, we can discuss the asymptotic behaviour and the dynamical system for the fractional differential equation. Here we will mention only few topics which should be exploited more.

Let us consider

$$D_t^\alpha U(t) = AU(t), \quad t > 0, \quad (3.4)$$

where $U = (u_1, \dots, u_N)^T$, \cdot^T denotes the transpose of the vector under consideration, and A is an $N \times N$ constant matrix. Then we can prove

Proposition 3.1

Let $0 < \alpha < 1$ and let the real parts of all the eigenvalues of A be negative. Then there exists a constant $C > 0$ such that

$$\|u(t)\| \leq \frac{C}{t^\alpha} \|u(0)\|, \quad t > 0$$

for an arbitrary solution to (3.4).

Unlike the case $\alpha = 1$, we can not have the exponential decay. Moreover the decay rate $t^{-\alpha}$ is the best possible as the following example shows: Let $N = 1$ and consider (3.2) with $\lambda < 0$. Then by Theorem 1.4 (pp.33-34) in [41] implies that the solution can not decay faster than $t^{-\alpha}$. In sections 4 and 5, we discuss similar properties of solutions of the fractional diffusion equation as $t \rightarrow \infty$.

In view of Proposition 3.1, we can discuss the linearized stability for

$$D_t^\alpha u(t) = Au + F(u, t)$$

with vector-valued function u and a suitable nonlinear term F .

As other interesting problem, we can mention the global existence in time to a nonlinear ordinary fractional differential equation. For example let us consider:

$$D_t^\alpha u(t) = -u(1 - u), \quad t > 0. \quad (3.5)$$

In the case of $\alpha = 1$, the following is well-known and can be proved easily.

$0 < u(0) < 1$: the solution exists globally.

$u(0) > 1$: the solution can not exist globally.

However for $0 < \alpha < 1$, such a result is not known. The difficulty comes from that $D_t^\alpha u(t)$ does not give information of $u(t)$ near t (the converse to Proposition 2.2 is not true).

We further mention a few works on the chaos for systems of ordinary fractional differential equations and refer to Ge and Hsu [11], Li and Peng [23] where chaoses are observed by numerical simulations for some systems. In the latter paper, the authors consider

$$\begin{cases} D_t^{\alpha_1} u = a(v - u), \\ D_t^{\alpha_2} v = (c - a)u - uv + cv, \\ D_t^{\alpha_3} w = uv - bw, \end{cases}$$

where $0 < \alpha_1, \alpha_2, \alpha_3 \leq 1$ and $a, b, c \in \mathbb{R}$. However comprehensive researches are not yet done. See also Luchko, Rivero, Trujillo and Pilar Velasco [27] which considers an inverse problem of determining a memory function in ordinary fractional differential equations.

4. FRACTIONAL DIFFUSION EQUATION

We survey results on the fractional diffusion equation (1.1).

The fractional diffusion equation has been introduced in physics by Nigmatullin [37] to describe diffusions in media with fractal geometry. One can regard (1.1) as a

macroscopic model derived from the continuous-time random walk. Metzler and Klafter [34] demonstrated that a fractional diffusion equation describes a non-Markovian diffusion process with a memory. See also Metzler, Glöckle and Nonnenmacher [32], Metzler and Klafter [33], Roman [43]. Roman and Alemany [44] investigated a continuous time random walks on fractals and showed that the average probability density of random walks on fractals obeys a diffusion equation with a fractional time derivative asymptotically. Ginoia, Cerbelli and Roman [13] presented a fractional diffusion equation describing relaxation phenomena in complex viscoelastic materials. Here we refer to several works on the mathematical treatments for equation (1.1). Kochubei [21], [22] applied the semigroup theory in Banach spaces, and Eidelman and Kochubei [8] constructed the fundamental solution in R^n and proved the maximum principle for the Cauchy problem. See also Mainardi [28] - [31] and Schneider and Wyss [49]. Gejji and Jafari [12] solved a nonhomogeneous fractional diffusion-wave equation in a 1-dimensional bounded domain. Fujita [10] discussed an integrodifferential equation which interpolates the heat equation and the wave equation in an unbounded domain. Agarwal [3] solved a fractional diffusion equation using a finite sine transform technique and presented numerical results in a 1-dimensional bounded domain.

We will solve equation (1.1) satisfying the following initial-boundary value conditions:

$$u(x, t) = 0, \quad x \in \partial\Omega, \quad t \in (0, T), \quad (5.1)$$

$$u(x, 0) = a(x), \quad x \in \Omega. \quad (5.2)$$

In spite of the importance, to the authors' best knowledge, there are not many works published concerning the unique existence of the solution to (1.1), (5.1) and (5.2) and the properties are remarkably different from the standard diffusion. In Prušs [42] (especially in Chapter I.3), one can refer to the methods for (1.1). See also Bazhlekova [4] and Gorenflo, Luchko and Zabrejko [14], Gorenflo and Mainardi [16].

The maximum principle for (1.1) with (5.1) is recently proved in Luchko [25] and see also a new paper Luchko [26] which proves the well-posedness of the forward problem (1.1), (5.1) and (5.2), but we will here present more detailed regularity and qualitative properties.

In particular, for discussions on inverse problems, we need representation formulae of the solution to (1.1), (5.1) and (5.2) by the eigenfunctions, and to the authors' best knowledge, there are no results published concerning the regularity properties of the eigenfunction expansions of the solutions which are corresponding to Chapter 3 of Lions and Magenes [24] and Pazy [39] for example.

In this section, we will show the well-posedness of the solution given by the Fourier method. Second we establish several uniqueness results for related inverse problems.

Let $L^2(\Omega)$ be a usual L^2 -space with the scalar product (\cdot, \cdot) . We denote the Sobolev spaces by $H^\ell(\Omega)$ with $\ell > 0$ (e.g., Adams [2]).

We define an operator L in $L^2(\Omega)$ by

$$(Lu)(x) = \sum_{i,j=1}^n \frac{\partial}{\partial x_i} \left(a_{ij}(x) \frac{\partial u}{\partial x_j}(x) \right) + c(x)u(x), \quad x \in \Omega,$$

$$\mathcal{D}(L) = H^2(\Omega) \cap H_0^1(\Omega).$$

Here and henceforth C_j denote positive constants which are independent of F in (1.1), a, b in (5.1) and (5.2), but may depend on α and the coefficients of the operator L .

Since $-L$ is a symmetric uniformly elliptic operator, the spectrum of L is entirely composed of eigenvalues and counting according to the multiplicities, we can set: $0 < \lambda_1 \leq \lambda_2 \leq \dots$. By $\varphi_n \in H^2(\Omega) \cap H_0^1(\Omega)$ we denote the orthonormal eigenfunction corresponding to $-\lambda_n$: $L\varphi_n = -\lambda_n\varphi_n$. Then the sequence $\{\varphi_n\}_{n \in \mathbb{N}}$ is orthonormal basis in $L^2(\Omega)$.

We are ready to state our main theorems on the unique existence of solution to (1.1), (5.1) and (5.2).

Theorem 4.1

Let $F = 0$.

(i) For $a \in L^2(\Omega)$, there exists a unique solution $u \in C([0, T]; L^2(\Omega)) \cap C((0, T]; H^2(\Omega) \cap H_0^1(\Omega))$ such that $\partial_t^\alpha u \in C((0, T]; L^2(\Omega))$. Moreover there exists a constant $C_1 > 0$ such that

$$\|u\|_{C([0, T]; L^2(\Omega))} \leq C_1 \|a\|_{L^2(\Omega)},$$

$$\|u(\cdot, t)\|_{H^2(\Omega)} + \|\partial_t^\alpha u(\cdot, t)\|_{L^2(\Omega)} \leq C_1 t^{-\alpha} \|a\|_{L^2(\Omega)}$$

for all $a \in L^2(\Omega)$. The eigenfunction expansion holds:

$$u(x, t) = \sum_{n=1}^{\infty} (a, \varphi_n) E_{\alpha, 1}(-\lambda_n t^\alpha) \varphi_n(x)$$

in $C([0, T]; L^2(\Omega)) \cap C((0, T]; H^2(\Omega) \cap H_0^1(\Omega))$.

(ii) There exists a constant $C_2 > 0$ such that

$$\|u\|_{L^2(0, T; H^2(\Omega))} + \|\partial_t^\alpha u\|_{L^2(\Omega \times (0, T))} \leq C_2 \|a\|_{H^1(\Omega)}$$

for all $a \in H_0^1(\Omega)$.

(iii) There exists a constant $C_3 > 0$ such that

$$\|u\|_{C([0, T]; H^2(\Omega))} + \|\partial_t^\alpha u\|_{C([0, T]; L^2(\Omega))} \leq C_3 \|a\|_{H^2(\Omega)}$$

for $a \in H^2(\Omega) \cap H_0^1(\Omega)$.

Theorem 4.2

Let $a = 0$ and $F \in L^2(\Omega \times (0, T))$. Then there exists a unique solution $u \in L^2(0, T; H^2(\Omega) \cap H_0^1(\Omega))$ and there exists a constant $C_4 > 0$ such that

$$\|u\|_{L^2(0, T; H^2(\Omega))} + \|\partial_t^\alpha u\|_{L^2(\Omega \times (0, T))} \leq C_4 \|F\|_{L^2(\Omega \times (0, T))}$$

for all $F \in L^2(\Omega \times (0, T))$. Moreover

$$u(x, t) = \sum_{n=1}^{\infty} \int_0^t (F(\cdot, t-\tau), \varphi_n) \tau^{\alpha-1} \times E_{\alpha, \alpha}(-\lambda_n \tau^\alpha) d\tau \varphi_n(x).$$

Remark. These results include the case of $\alpha = 1$

For $\theta \in (0, 1)$, we set

$$\|F\|_{C^\theta([0, T]; L^2(\Omega))} = \|F\|_{C([0, T]; L^2(\Omega))} + \sup_{0 \leq t < s \leq T} \frac{\|F(\cdot, t) - F(\cdot, s)\|_{L^2(\Omega)}}{|t - s|^\theta}.$$

Next we show the maximal regularity for $F \in C^\theta([0, T]; L^2(\Omega))$.

Theorem 4.3

Let $a \in L^2(\Omega)$ and $F \in C^\theta([0, T]; L^2(\Omega))$.

The solution u is represented by eigenfunction expansion:

$$u(x, t) = \sum_{n=1}^{\infty} \left\{ (a, \varphi_n) E_{\alpha, 1}(-\lambda_n t^\alpha) + \int_0^t (F(\cdot, t-\tau), \varphi_n) \tau^{\alpha-1} E_{\alpha, \alpha}(-\lambda_n \tau^\alpha) d\tau \right\} \varphi_n(x).$$

(1) For arbitrary $\delta > 0$, there exists a constant $C_5 = C_5(\delta) > 0$ such that

$$\|Lu\|_{C^\theta([\delta, T]; L^2(\Omega))} + \|\partial_t^\alpha u\|_{C^\theta([\delta, T]; L^2(\Omega))} \leq \frac{C_5}{\delta} (\|F\|_{C^\theta([\delta, T]; L^2(\Omega))} + \|a\|_{L^2(\Omega)})$$

for all $a \in L^2(\Omega)$ and $F \in C^\theta([0, T]; L^2(\Omega))$.

(2) There exists a constant $C_6 > 0$ such that

$$\|Lu\|_{C([0, T]; L^2(\Omega))} + \|\partial_t^\alpha u\|_{C([0, T]; L^2(\Omega))} \leq C_6 (\|a\|_{H^2(\Omega)} + \|F\|_{C^\theta([0, T]; L^2(\Omega))})$$

for $a \in H^2(\Omega) \cap H_0^1(\Omega)$ and $F \in C^\theta([0, T]; L^2(\Omega))$.

(3) Let $a = 0$. There exists a constant $C_7 > 0$ such that

$$\|Lu\|_{C^\theta([0, T]; L^2(\Omega))} + \|\partial_t^\alpha u\|_{C^\theta([0, T]; L^2(\Omega))} \leq C_7 \|F\|_{C^\theta([0, T]; L^2(\Omega))}$$

for all $F \in C^\theta([0, T]; L^2(\Omega))$ satisfying $F(\cdot, 0) = 0$.

This is the same as the case of $\alpha = 1$. Prüss [42] already proved Theorem 4.3 (3).

Corollary 4.1 (slow decay)

Let $a \in L^2(\Omega)$ and $F = 0$. Then there exist constants $C_8, C_9 > 0$ such that

$$\|u(\cdot, t)\|_{L^2(\Omega)} \leq \frac{C_8}{1 + \lambda_1 t^\alpha} \|a\|_{L^2(\Omega)}, \quad t \geq 0,$$

$$\|\partial_t^m u(\cdot, t)\|_{L^2(\Omega)} \leq \frac{C_9}{t^m} \|a\|_{L^2(\Omega)}, \quad t > 0, m \in \mathbb{N}$$

for all $a \in L^2(\Omega)$. Here C_9 is independent of m .

Here we compare our results with the case $\alpha = 1$: $t^{-\alpha}$ -decay for $0 < \alpha < 1$ but the exponential decay for $\alpha = 1$.

We can consider $\alpha > 1$ similarly and see Sakamoto [45], Sakamoto and Yamamoto [46].

5. FURTHER QUALITATIVE RESULTS FOR THE FRACTIONAL DIFFUSION EQUATION

5.1. BACKWARD PROBLEM IN TIME

It is well-known that the backward problem in time is severely ill-posed for the parabolic problem (i.e., $\alpha = 1$). The severe ill-posedness means that we can not recover the stability in the backward problem even if we strengthen the norm within Sobolev norms for estimating the initial value in $L^2(\Omega)$. For $0 < \alpha < 1$, the backward problem is moderately ill-posed, as the following theorem implies:

Theorem 5.1

Let $0 < \alpha < 1$. For arbitrary $T > 0$ and arbitrary $a_1 \in H^2(\Omega) \cap H_0^1(\Omega)$, there exists a unique solution $u \in C([0, T]; L^2(\Omega)) \cap C((0, T]; H^2(\Omega) \cap H_0^1(\Omega))$ such that $u(\cdot, T) = a_1$ to the forward problem (1.1), (5.1) and (5.2) with $F = 0$. Moreover there exist constants $C_{10}, C_{11} > 0$ such that

$$\begin{aligned} C_{10} \|u(\cdot, 0)\|_{L^2(\Omega)} &\leq \|u(\cdot, T)\|_{H^2(\Omega)} \\ &\leq C_{11} \|u(\cdot, 0)\|_{L^2(\Omega)}. \end{aligned}$$

5.2. UNIQUENESS OF SOLUTION

The solution can be uniquely determined by data in any small subdomain over time interval. This is closely related with the approximate controllability (e.g., Georg Schmidt and Weck [48]) but we will omit further discussions.

Theorem 5.2

Let $0 < \alpha < 1$.

Let spatial dimension ≤ 3 , $a \in H_0^4(\Omega)$,

$$\partial_t^\alpha u = \sum_{i,j=1}^n \frac{\partial}{\partial x_i} (a_{ij}(x) \frac{\partial}{\partial x_j} u) + c(x)u,$$

$u|_{\partial\Omega} = 0$ and $u = 0$ in $\omega \times (0, T)$ with arbitrary subdomain ω and $T > 0$. Then $u = 0$ in $\Omega \times (0, T)$.

5.3. DECAY AT $t = \infty$

Non-trivial solutions can not decay faster than polynomial orders, which implies the slow diffusion for $0 < \alpha < 1$. See also Corollary 4.1 in section 4.

Theorem 5.3 Let $0 < \alpha < 1$, ω be an arbitrary subdomain, let spatial dimension ≤ 3 and $a \in H_0^4(\Omega)$, $\partial_t^\alpha u = \sum_{i,j=1}^n \frac{\partial}{\partial x_i} (a_{ij}(x) \frac{\partial}{\partial x_j} u) + c(x)u$, $u|_{\partial\Omega} = 0$. Let for all $m \in \mathbb{N}$, there exists a constant $C(m) > 0$ such that $\|u(\cdot, t)\|_{L^\infty(\omega)} \leq \frac{C(m)}{t^m}$ as $t \rightarrow \infty$. Then $u = 0$ in $\Omega \times (0, \infty)$.

5.4. OTHER INVERSE PROBLEM

Let $p > 0$ on $[0, \ell]$ and $p \in C^2([0, \ell])$. We consider

$$\partial_t^\alpha u(x, t) = \frac{\partial}{\partial x} \left(p(x) \frac{\partial u}{\partial x} \right), \quad 0 < x < \ell, \quad 0 < t < T,$$

$$u(x, 0) = \delta(x) : \text{delta function,}$$

$$u_x(0, t) = u_x(\ell, t) = 0.$$

It is practically difficult to determine the order α a priori and it is important to determine the order α and the diffusion coefficient $p(x)$ by available observation data at the boundary point $x = 0$ over time interval. Thus the following inverse problem is significant.

Inverse problem: Determine $\alpha \in (0, 1)$ and $p(x)$, $0 < x < \ell$ by $u(0, t)$, $0 < t < T$.

Then the uniqueness is proved in Cheng, Nakagawa, Yamamoto and Yamazaki [7] by means of the Gel'fanf-Levitan theory (see e.g., Freiling and Yurko [9]) and the eigenfunction expansion. Fixed $\alpha = 1$, a similar inverse problem is considered in Murayama [36], Pierce [40], Suzuki and Murayama [51]. By the results in section 4, we can consider other types of inverse problems and we refer to Isakov [19] as monographs on inverse problems for partial differential equations.

ACKNOWLEDGEMENTS

The second named author was supported partly by the 21st Century COE program, the Global COE program and the Doctoral Course Research Accomplishment Cooperation System at Graduate School of Mathematical Sciences of The University of Tokyo, and the Japan Student Services Organization.

REFERENCES

- [1] E.E. Adams and L.W. Gelhar, Field study of dispersion in a heterogeneous aquifer 2. spatial moments analysis, *Water Resources Research* **28** (1992) 3293–3307.
- [2] R. A. Adams, *Sobolev Spaces*, Academic Press, New York, 1975.
- [3] O. P. Agarwal, Solution for a fractional diffusion-wave equation defined in a bounded domain, *Nonlinear Dynamics* **29** (2002) 145–155.
- [4] E. Bazhlekova, The abstract Cauchy problem for the fractional evolution equation, *Fractional Calculus & Applied Analysis* **1** (1998) 255–270.
- [5] B. Berkowitz, A. Cortis, M. Dentz and H. Scher, Modeling non-Fickian transport in geological formations as a continuous time random walk, *Review of Geophysics* **44** (2006), RG2003/2006, Paper number 2005RG000178.

- [6] B. Berkowitz, H. Scher and S.E. Silliman, Anomalous transport in laboratory-scale, heterogeneous porous media, *Water Resources Research* **36** (2000) 149–158.
- [7] J. Cheng, J. Nakagawa, M. Yamamoto and T. Yamazaki, Uniqueness in an inverse problem for one-dimensional fractional diffusion equation, *Inverse Problems* **25** (2009) 115002 (16 pp).
- [8] S.D. Eidelman and A.N. Kochubei, Cauchy problem for fractional diffusion equations, *J. Differential Equations* **199** (2004) 211–255.
- [9] G. Freiling and V. Yurko, *Inverse Sturm-Liouville Problems and Their Applications*, Nova Science Publ. New York, 2001.
- [10] Y. Fujita, Integrodifferential equation which interpolates the heat equation and the wave equation, *Osaka J. Math.* **27** (1990), 309–321, 797–804.
- [11] Z.-M. Ge and M.-Y. Hsu, Chaos excited chaos synchronization of integral and fractional order generalized van der Pol systems, *Chaos, Solitons and Fractals* **36** (2008) 592–604.
- [12] V. D. Gejji and H. Jafari, Boundary value problems for fractional diffusion-wave equation, *Aust. J. Math. Anal. and Appl.* **3** (2006), 1–8.
- [13] M. Ginoa, S. Cerbelli and H. E. Roman, Fractional diffusion equation and relaxation in complex viscoelastic materials, *Physica A* **191** (1992), 449–453.
- [14] R. Gorenflo, Yu. F. Luchko and P.P. Zabrejko, On solvability of linear fractional differential equations in Banach spaces, *Fractional Calculus & Applied Analysis* **2** (1999) 163–176.
- [15] R. Gorenflo and F. Mainardi, Fractional oscillations and Mittag-Leffler functions, *International Workshop on the Recent Advances in Applied Mathematics (Kuwait, RAAM'96)*, Kuwait, 1996, 193–208.
- [16] R. Gorenflo and F. Mainardi, Fractional calculus: Integral and differential equations of fractional order, in "Fractals and Fractional Calculus in Continuum Mechanics", Springer-Verlag, New York, 1997. 223–276.
- [17] R. Gorenflo and R. Rutman, On ultraslow and intermediate processes, in "International Workshop on Transforms Methods and Special Functions", Bulgarian Acad. Sci. Sofia, 1994, 61–81.
- [18] Y. Hatano and N. Hatano, Dispersive transport of ions in column experiments: an explanation of long-tailed profiles, *Water Resources Research* **34** (1998) 1027–1033.
- [19] V. Isakov, *Inverse Problems for Partial Differential Equations*, Springer-Verlag, Berlin, 2006.
- [20] A. A. Kilbas, H. M. Srivastava and J. J. Trujillo, *Theory and Applications of Fractional Differential Equations*, Elsevier, Amsterdam, 2006.
- [21] A. N. Kochubei, A Cauchy problem for evolution equations of fractional order, *J. Diff. Equ.* **25** (1989) 967–974.
- [22] A. N. Kochubei, Fractional order diffusion, *J. Diff. Equ.* **26** (1990) 485–492.
- [23] C. Li and G. Peng, Chaos in Chen's system with a fractional order, *Chaos, Solitons and Fractals* **22** (2004) 443–450.
- [24] J.-L. Lions and E. Magenes, *Non-homogeneous Boundary Value Problems and Applications*, Vol.I, II, Springer-Verlag, Berlin, 1972.
- [25] Y. Luchko, Maximum principle for the generalized time-fractional diffusion equation, *J. Math. Anal. Appl.* **351** (2009) 218–223.
- [26] Y. Luchko, Some uniqueness and existence results for the initial-boundary-value problems for the generalized time-fractional diffusion equation, *Computers and Mathematics with Applications* (2009), doi: 10.1016/j.camwa.2009.08.015
- [27] Y. Luchko, M. Rivero, J.J. Trujillo and M. Pilar Velasco, Fractional models, non-locality, and complex systems, *Computers and Mathematics with Applications* (2009), doi: 10.1016/j.camwa.2009.05.018
- [28] F. Mainardi, On the initial value problem for the fractional diffusion-wave equation, in "Waves and Stability in Continuous Media", World Scientific, Singapore, 1994, 246–251.
- [29] F. Mainardi, The time fractional diffusion-wave equation, *Radiophys. and Quant. Elect.* **38** (1995) 13–24.
- [30] F. Mainardi, The fundamental solutions for the fractional diffusion-wave equation, *Appl. Math. Lett.* **9** (1996) 23–28.
- [31] F. Mainardi, Fractional calculus: some basic problems in continuum and statistical mechanics, in "Fractals and Fractional Calculus in Continuum Mechanics", Springer-Verlag, New York, 1997, 291–348.
- [32] R. Metzler, W.G. Glöckle and T.F. Nonnenmacher, Fractional model equation for anomalous diffusion, *Physica A* **211** (1994) 13–24.
- [33] R. Metzler and J. Klafter, Boundary value problems for fractional diffusion equations, *Physica A* **278** (2000) 107–125.
- [34] R. Metzler and J. Klafter, The random walk's guide to anomalous diffusion: a fractional dynamics approach, *Physics Reports* **339** (2000) 1–77.

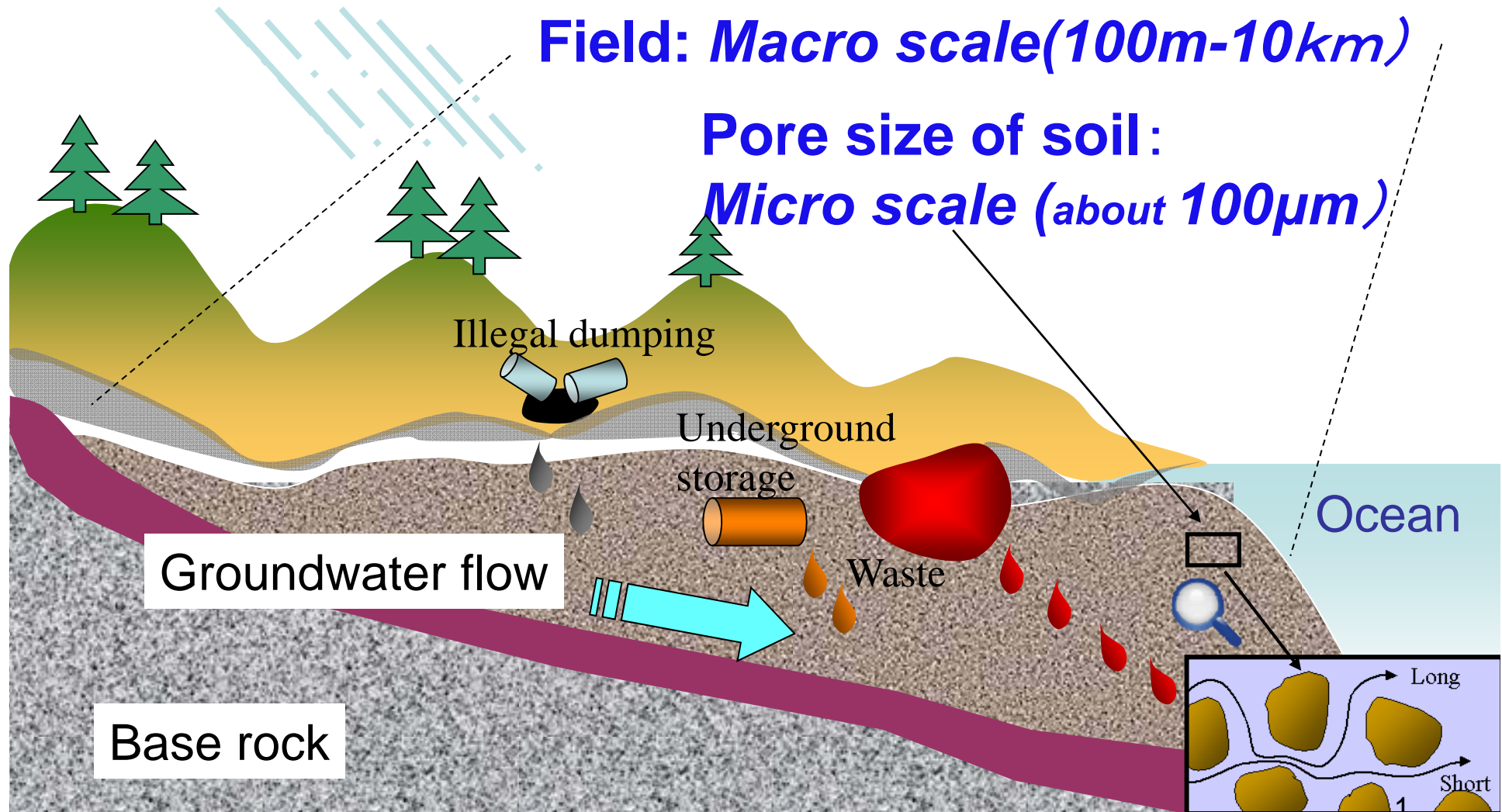
- [35] K. S. Miller and B. Ross, *An Introduction to the Fractional Calculus and Fractional Differential Equations*, Wiley, New York, 1993.
- [36] R. Murayama, The Gel'fand-Levitan theory and certain inverse problems for the parabolic equation, *J. Fac. Sci. The Univ. Tokyo, Section IA, Math.* **28** (1981) 317–330.
- [37] R. R. Nigmatullin, The realization of the generalized transfer equation in a medium with fractal geometry, *Phys. Stat. Sol. B* **133** (1986) 425–430.
- [38] K. B. Oldham and J. Spanier, *The Fractional Calculus: Theory and Applications of Differentiation and Integration to Arbitrary Order*, Academic Press, New York, 1974.
- [39] A. Pazy, *Semigroups of Linear Operators and Applications to Partial Differential Equations*, Springer-Verlag, Berlin, 1983.
- [40] A. Pierce, Unique identification of eigenvalues and coefficients in a parabolic problem, *SIAM J. Control and Optim.* **17** (1979) 494–499.
- [41] I. Podlubny, *Fractional Differential Equations*, Academic Press, San Diego, 1999.
- [42] J. Prüss, *Evolutionary Integral Equations and Applications*, Birkhäuser, Basel, 1993.
- [43] H.E. Roman, Structures of random fractals and the probability distribution of random walks, *Phys. Rev. E* **51** (1995) 5422–5425.
- [44] H. E. Roman and P. A. Alemany, Continuous-time random walks and the fractional diffusion equation, *J. Phys. A* **27** (1994) 3407–3410.
- [45] K. Sakamoto, *Inverse Source Problems for Diffusion Equations and Fractional Diffusion Equations*, Ph.D. thesis, Graduate School of Mathematical Sciences, The University of Tokyo, 2009.
- [46] K. Sakamoto and M. Yamamoto, Initial value/boundary value problems for fractional diffusion-wave equations and applications to some inverse problems, preprint.
- [47] S. G. Samko, A. A. Kilbas and O. I. Marichev, *Fractional Integrals and Derivatives*, Gordon and Breach Science Publishers, Philadelphia, 1993.
- [48] E.J.P. Georg Schmidt and N. Weck, On the boundary behavior of solutions to elliptic and parabolic equations - with applications to boundary control for parabolic equations, *SIAM J. Control and Optim.* **16** (1978) 593–598.
- [49] W. R. Schneider and W. Wyss, Fractional diffusion and wave equations, *J. MathPhys.* **30** (1989) 134–144.
- [50] I.M. Sokolov, J. Klafter and A. Blumen, Fractional kinetics, *Physics Today* **55** (2002) 48–54.
- [51] T. Suzuki and R. Murayama, A uniqueness theorem in an identification problem for coefficients of parabolic equations, *Proc. Japan Acad. Ser. A* **56** (1980) 259–263.
- [52] Y. Xiong, G. Huang and Q. Huang, Modeling solute transport in one-dimensional homogeneous and heterogeneous soil columns with continuous time random walk, *J. Contaminant Hydrology* **86** (2006) 163–175.

Junichi Nakagawa, Kenichi Sakamoto
 Mathematical Science & Technology Research Lab., Advanced Technology Research Laboratories, Technical Development Bureau, Nippon Steel Corporation, 20-1 Shintomi Futtsu Chiba, 293-8511, Japan
 E-mail: nakagawa.junichi(at)nsc.co.jp
 E-mail: kens(at)mf.nacsinet.com

Masahiro Yamamoto
 Department of Mathematical Sciences, The University of Tokyo, Komaba Meguro Tokyo 153-8914 Japan
 E-mail: myama(at)ms.u-tokyo.ac.jp

Issues Seen by Academia Engineering Researchers

“The Prediction of the Progress of Soil Contamination”



Model Prediction and Reality

